

# Orifice Flow Resistance

Apply Bernoulli's equation for streamlines between 1 and 2, 2 and 3,

$$\frac{P_1}{\rho} + \frac{v_1^2}{2} = \frac{P_2}{\rho} + \frac{v_2^2}{2} = \frac{P_3}{\rho} + \frac{v_3^2}{2}$$

Equating volume flow rate at all cross sections,

$$Q = A_1 v_1 = A_2 v_2 = A_3 v_3 = A_4 v_4$$

Assume sudden expansion between 2 and 4, then there is a velocity head loss and pressures are same at 3 and 4, hence  $P_3 = P_4$  (easy to measure)

Assume  $A_1 > A_2$  and  $A_1 > A_3$ , then  $v_1 < v_3$  and  $v_4 < v_3$ , so  $v_1$  and  $v_4$  are negligible

Combine above equations with assumptions,

$$P_1 - P_4 = \frac{\rho}{2} v_3^2$$

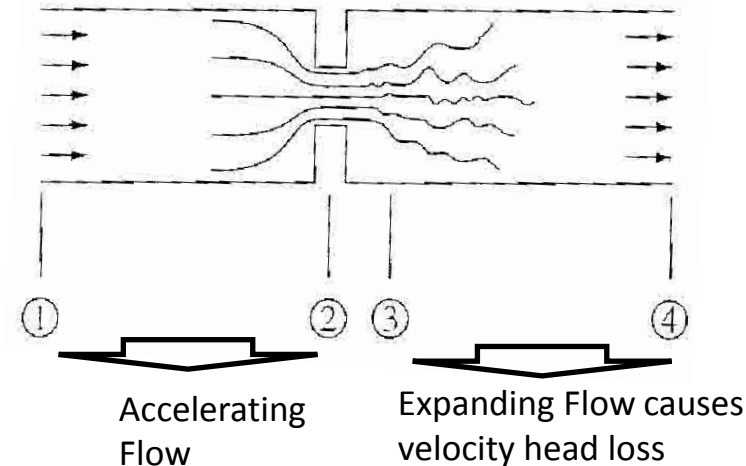
Use discharge coefficient  $C_d$  (combined effects of losses between 1 and 3, and representing  $A_3$  in term of  $A_2$ )

$$\delta P = P_1 - P_4 = \frac{\rho}{2 C_d^2 A_2^2} Q^2 \quad (\text{orifice flow relation})$$

$$\text{or } Q = C_d A_2 \sqrt{\frac{2 \delta P}{\rho}} = \sqrt{\frac{|\delta P|}{K}} \text{sign}(\delta P)$$

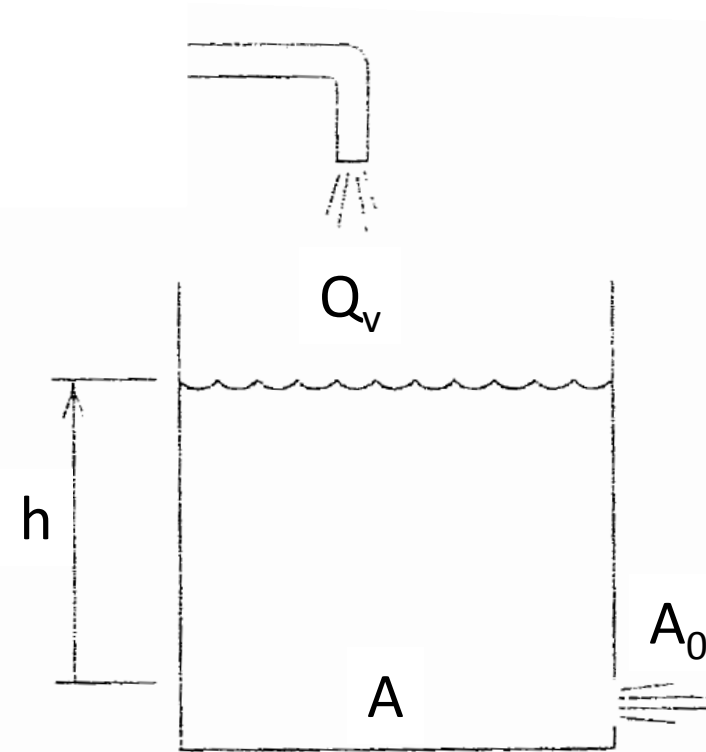
$$\text{where } K = \frac{\rho}{2 C_d^2 A_2^2}; \text{sign}(x) = 1 \text{ if } x > 0; \text{sign}(x) = -1 \text{ if } x < 0$$

$$C_d = C_d(\text{geometry, roughness, orifice configuration}, N_r) \sim .5 \text{ to } 1$$



# Example 5: Tank with an Orifice

- The tank shown has an orifice in its side wall. The orifice area is  $A_0$  and the bottom area of the tank is  $A$ . The liquid height above the orifice  $h$ . The volume inflow rate is  $Q_v$ . Develop a model of the height  $h$  with  $Q_v$  as the input.



# Example 5: Tank with an Orifice

From conservation of mass and the orifice flow relation,

$$\dot{m}_{\text{tank}} = \dot{m}_{in} - \dot{m}_{out}$$

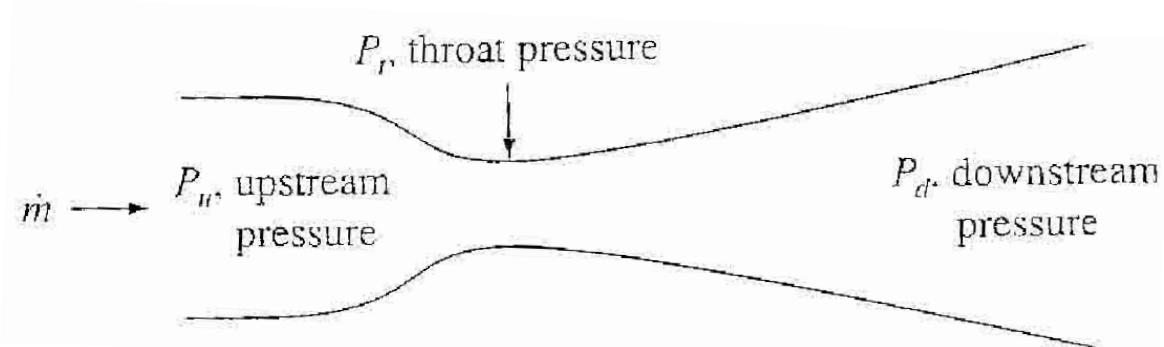
$$\rho A \frac{dh}{dt} = \rho Q_v - \rho C_d A_0 \sqrt{\frac{2\delta P}{\rho}}$$

$$\text{but } \delta P = \rho gh,$$

$$\rho A \frac{dh}{dt} = \rho Q_v - \rho C_d A_0 \sqrt{2gh}$$

This is a nonlinear equation, how can it be analyzed?

# Compressible Flow Resistance



- Similar to orifice flow, but considers density variation and has higher degree of nonlinearity and complex equations
- For compressible flow, mach number,  $N_m = v/c_0$ ,
  - cannot exceed 1 at the throat
  - equal 1 for air at  $P_t/P_u = .528$  (flow is choked)
- Once choked,  $N_m$  at throat stays at 1  $\rightarrow Q = A_t c_0 = \text{constant}$
- Due to compressibility, density can vary  $\rightarrow$  mass flow rate ( $=\rho_t Q$ ) can vary if upstream conditions change

# Compressible Flow Resistance

The critical pressure ratio determines whether flow is choked,

$$P_{cr} = \left( \frac{2}{k+1} \right)^{\frac{k}{k-1}}$$

*Mass flow rate per unit area,*

$$\frac{\dot{m}}{C_d A} = \sqrt{\frac{2}{RT^*}} \sqrt{\frac{T^*}{T_u}} \sqrt{\frac{k}{k-1}} P_u \left[ \left( \frac{P_t}{P_u} \right)^{\frac{2}{k}} - \left( \frac{P_t}{P_u} \right)^{\frac{k+1}{k}} \right]^{\frac{1}{2}} \quad (\text{unchoked flow, } \frac{P_t}{P_u} > P_{cr})$$

$$\frac{\dot{m}}{C_d A} = \sqrt{\frac{2}{RT^*}} \sqrt{\frac{T^*}{T_u}} \sqrt{\frac{k}{k-1}} P_u \left[ (P_{cr})^{\frac{2}{k}} - (P_{cr})^{\frac{k+1}{k}} \right]^{\frac{1}{2}} \quad (\text{choked flow, } \frac{P_t}{P_u} = P_{cr})$$

where  $T^* = \text{reference temperature}$

$P_t$  is not known usually. For orifice exits into large chamber, assume  $P_t = P_d$ .

One disadvantage of compressible equations : ability to solve for mass flow rate in terms of pressures, but not reverse.

# Compressible Flow Resistance

- Approximation to compressible flow equation is possible by using incompressible flow equation  $\delta P = \rho Q^2 / (2C_d^2 A^2)$  and gas law, and substituting throat pressure with upstream pressure.

$$\frac{\dot{m}}{C_d A} = \sqrt{\frac{2}{RT^*}} \sqrt{\frac{T^*}{T_u}} P_u \left[ \left( \frac{P_t}{P_u} \right) - \left( \frac{P_t}{P_u} \right)^2 \right]^{\frac{1}{2}} \quad (\text{unchoked incompressible approx})$$

$$\frac{\dot{m}}{C_d A} = \sqrt{\frac{2}{RT^*}} \sqrt{\frac{T^*}{T_u}} \sqrt{\frac{k}{k-1}} P_u \left[ (P_{cr}) - (P_{cr})^2 \right]^{\frac{1}{2}} \quad (\text{choked incompressible approx})$$

- Can solve for  $P_t$  or  $P_u$  as a function of mass flow rate

# Case Study: Spring-Loaded Diaphragm Actuator

# Spring-Loaded Diaphragm Actuator

- Read handout
- Work in groups to answer parts a, b, and c



# Spring-Loaded Diaphragm Actuator

Model of system using incompressible orifice equation and no mass:

$$\delta P_0 - \delta P_1 = K_o Q^2 \text{sign}(Q), \quad K_o = \frac{\rho}{2 C_d^2 A_o^2}, \quad A \delta P_1 - k_s z = 0$$

$$Q = \frac{V}{\beta} \delta \dot{P}_1 + \dot{V}, \quad V = V_o + A z, \quad \dot{V} = A \dot{z}, \quad \beta = n(\delta P_1 + P_{atm})$$

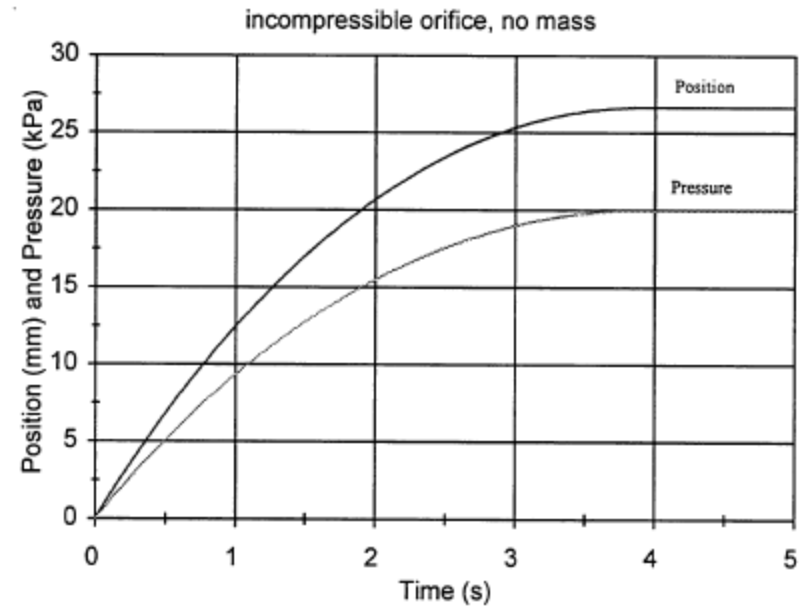
$$\dot{z} = \frac{\sqrt{\left| \delta P_0 - \frac{k_s}{A} z \right|} \text{sign}\left(\delta P_0 - \frac{k_s}{A} z\right)}{\left[ \frac{V}{\beta} \frac{k_s}{A^2} + 1 \right]}$$

state variables:  $u_1 = \delta P_0, \quad x_1 = z$

# Spring-Loaded Diaphragm Actuator

$$\dot{x}_1 = \frac{\sqrt{\frac{u_1 - \frac{k_s}{A} x_1}{K_o A^2}} \operatorname{sign}(u_1 - \frac{k_s}{A} x_1)}{\left[ \frac{(V_o + A x_1)}{n \left( \frac{k_s}{A} x_1 + P_{atm} \right)} \frac{k_s}{A^2} + 1 \right]}$$

outputs:  $y_1 = z = x_1$ ,  $y_2 = \delta P_1 = \frac{k_s}{A} x_1$



# Spring-Loaded Diaphragm Actuator

Model of system using incompressible orifice equation with mass:

Change force balance to  $A \delta P_1 - M \ddot{z} - k_s z = 0$ :

$$\left[ M D^2 + k_s \right] z = A \delta P_1, \quad \frac{V}{\beta} \delta \dot{P}_1 = \sqrt{\frac{|\delta P_0 - \delta P_1|}{K_o}} \text{sign}(\delta P_0 - \delta P) - A \dot{z}$$

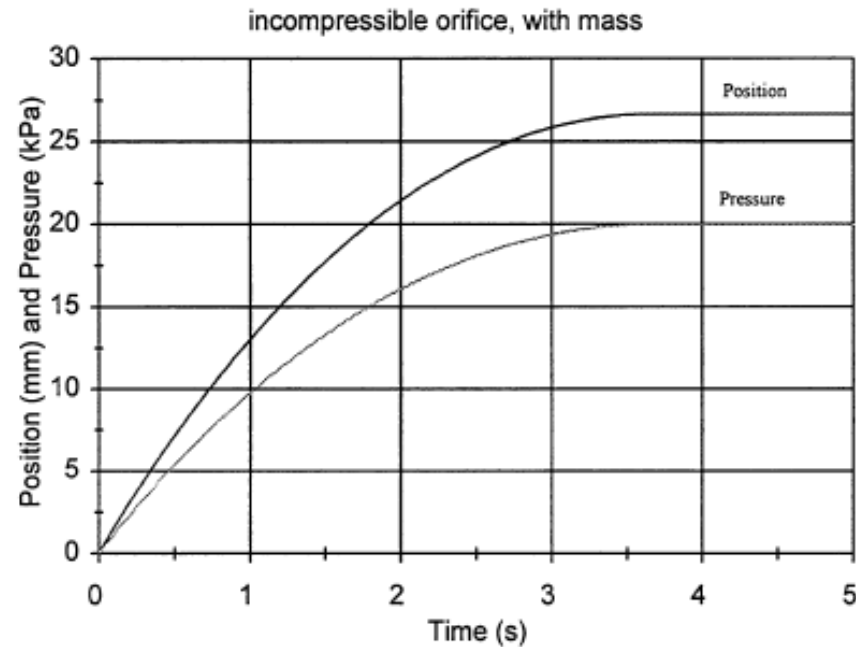
state variables:  $u_1 = \delta P_0$ ,  $x_1 = z$ ,  $\dot{x}_1 = \dot{z}$ ,  $x_2 = \delta P_1$

$$\dot{x}_1 = x_2, \quad \dot{x}_2 = -\frac{k_s}{M} x_1 + \frac{A}{M} x_2$$

$$\dot{x}_3 = \frac{n(x_3 + P_{atm})}{(V_o + A x_1)} \left\{ \sqrt{\frac{|u_1 - x_3|}{K_o}} \text{sign}(u_1 - x_3) - A x_2 \right\}$$

outputs:  $y_1 = z = x_1$ ,  $y_2 = \delta P_1 = x_2$

# Spring-Loaded Diaphragm Actuator



# Spring-Loaded Diaphragm Actuator

Model of system using compressible orifice equation (incompressible approximation) with mass:

Change flow equation to (neglecting temperature variations):

$$\dot{m}_{in} = C_d A_o \sqrt{\frac{2}{R T^*}} (\delta P_0 + P_{atm}) \sqrt{|P_r - P_r^2|} \text{sign}(P_r - P_r^2) , \quad C_d A_o = \sqrt{\frac{\rho_o}{2 K_o}}$$

where  $P_r = \frac{\delta P_1 + P_{atm}}{\delta P_0 + P_{atm}}$  and if  $P_r \leq 0.528$  , then  $P_r = 0.528$

Change continuity equation to:

$$\delta \dot{P}_1 = \frac{\beta}{V} \left[ \frac{\dot{m}}{\rho_1} - A \dot{z} \right] , \quad \rho_1 = \frac{\delta P_1}{R T_1}$$

state variables:  $u_1 = \delta P_0$  ,  $x_1 = z$  ,  $\dot{x}_1 = \dot{z}$  ,  $x_2 = \delta P_1$

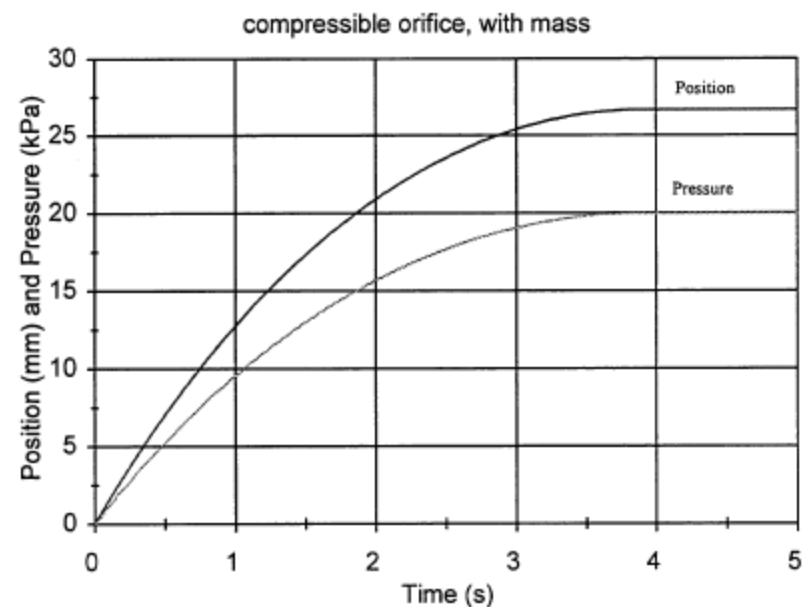
# Spring-Loaded Diaphragm Actuator

$$\dot{x}_1 = x_2, \quad \dot{x}_2 = -\frac{k_s}{M} x_1 + \frac{A}{M} x_2$$

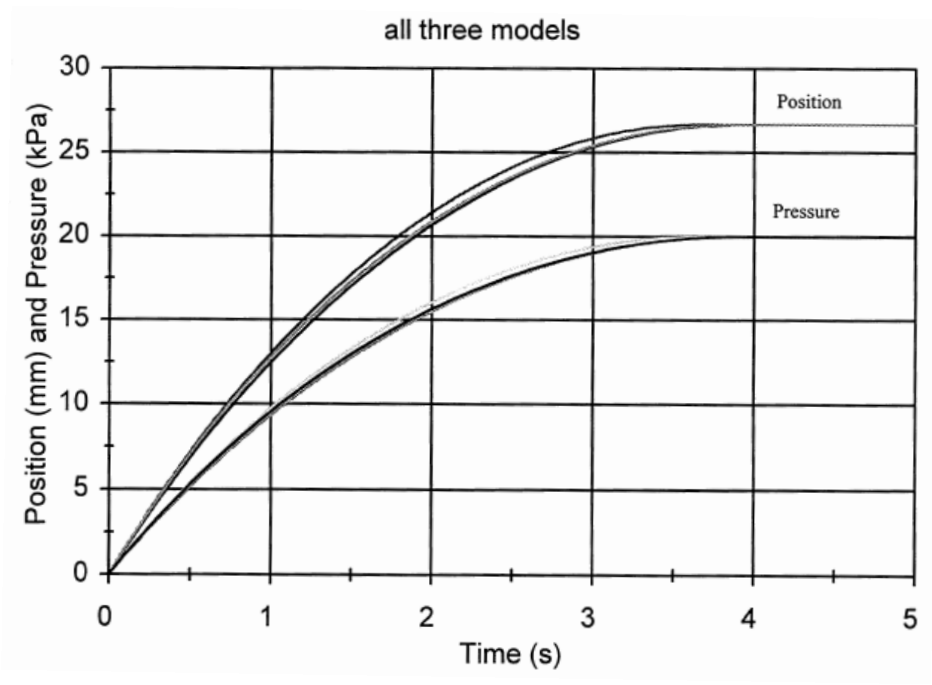
$$P_r = \frac{x_3 + P_{atm}}{u_1 + P_{atm}} \text{ and if } P_r \leq 0.528, \text{ then } P_r = 0.528$$

$$\dot{x}_3 = \frac{n(x_3 + P_{atm})}{(V_o + A x_1)} \left\{ C_d A_o \sqrt{2 R T_1} \frac{(\delta P_0 + P_{atm})}{(\delta P_1 + P_{atm})} \sqrt{|P_r - P_r^2|} \text{sign}(P_r - P_r^2) - A x_2 \right\}$$

outputs:  $y_1 = z = x_1, \quad y_2 = \delta P_1 = x_3$

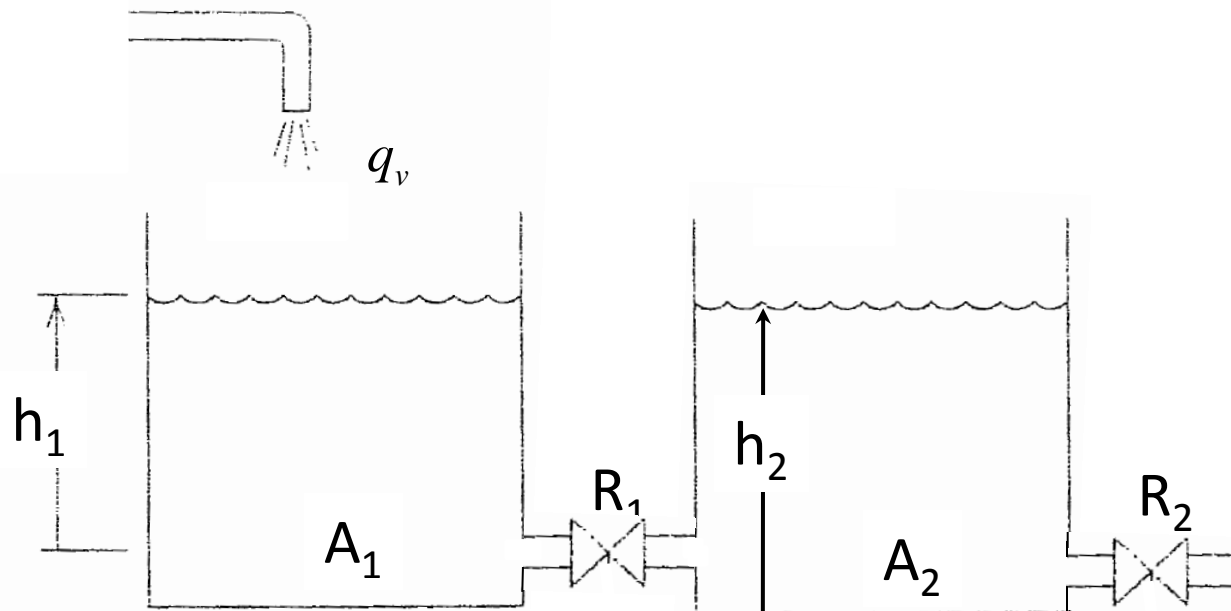


# Spring-Loaded Diaphragm Actuator



# Homework 5: Chapter 5

- 5.5, 5.11, 5.13, 5.15
- Two Interconnected Tanks Problem: Develop a model for the heights  $h_1$  and  $h_2$  in the liquid system shown in Figure below. The input volume flow rate  $q_v$  is given. Assume that laminar flow exists in the pipes. The laminar resistances are  $R_1$  and  $R_2$ , and the bottom areas of the tanks are  $A_1$  and  $A_2$ .





# References

- Woods, R. L., and Lawrence, K., Modeling and Simulation of Dynamic Systems, Prentice Hall, 1997.
- Close, C. M., Frederick, D. H., Newell, J. C., Modeling and Analysis of Dynamic Systems, Third Edition, Wiley, 2002
- Palm, W. J., Modeling, Analysis, and Control of Dynamic Systems